

Sequence of Maps Between Hopf and Aharonov-Bohm Bundles

M. Socolovsky

Instituto de Ciencias Nucleares, Universidad Nacional Autónoma de México, Circuito Exterior, Ciudad Universitaria, 04510, México D. F., México
 Email: socolovs@nucleares.unam.mx

Abstract The existence of the Aharonov-Bohm ($A - B$) effect with its associated $U(1)$ -bundle ξ_{A-B} , and the uniqueness -up to homotopy- of a continuous function $S^2 \rightarrow \mathbb{C}^*$, induce a unique map -up to isomorphism- between the Hopf bundles with zero and unit Chern number, respectively $\xi_0 : S^1 \rightarrow S^2 \times S^1 \rightarrow S^2$ and $\xi_1 : S^1 \rightarrow S^3 \rightarrow S^2$. This establishes a tight relation between ξ_0 and ξ_1 through ξ_{A-B} , and therefore between the $A - B$ effect and the hypothetical unit magnetic charge when the Dirac connection in ξ_1 is considered.

Keywords: Aharonov-Bohm effect, Hopf bundles, magnetic charge

1 Introduction

In a recent paper [1] we showed that, if $\xi_{A-B} : S^1 \rightarrow (T_0^2)^* \xrightarrow{pr_1} (D_0^2)^*$ and $\xi_D \equiv \xi_1 : S^1 \rightarrow S^3 \xrightarrow{\pi_D} S^2$ are the $U(1)$ -bundles where the Aharonov-Bohm ($A - B$) [2] effect and the unit Dirac monopole (D) [3,4] are represented, then ξ_{A-B} turns out to be isomorphic to the pull-back of ξ_D induced by the inclusion of the corresponding base spaces: $\iota : \mathbb{C}^* \rightarrow S^2$. ($S^1 = U(1)$, $S^2 \cong \mathbb{C} \cup \{\infty\}$, $(D_0^2)^* \cong \mathbb{R}^{2*} \cong \mathbb{C}^*$ is the punctured open 2-disk, $(T_0^2)^* = (D_0^2)^* \times S^1$ is the open solid 2-torus minus a circle, S^k , $k=1,2,3$, are spheres, pr_1 is the projection in the first entry, and π_D is the Hopf map $\pi_D(z_1, z_2) = z_1/z_2$ if $z_2 \neq 0$ and ∞ if $z_2 = 0$, with $\|(z_1, z_2)\| = 1$.) In particular, in the context of bundle theory, this allows to prove through the pull-back of the Dirac connection and the use of the Dirac quantization condition, the well known fact that the $A - B$ effect vanishes when the flux in the solenoid equals an integer times the quantum of magnetic flux $2\pi\hbar c/|e|$ associated with the electric charge $|e|$. Thus, one obtains a consequence on the observed $A - B$ effect [5,6] from the up to now only hypothetical existence of magnetic charges i.e. a sort of "passage" from D to $A - B$.

An immediate question that arises is if it is possible an "inverse passage": from $A - B$ to D , which might, at least from a purely mathematical point of view, reinforce the idea of the real existence of magnetic charges. In this note, we try to come close to a positive answer to this question using the fact that, up to homotopy, the *unique* continuous map from the base space of ξ_D to the base space of ξ_{A-B} is a constant map:

$$\kappa : S^2 \rightarrow \mathbb{C}^*, \quad b \mapsto \kappa(b) = z_0 \tag{1}$$

where z_0 is an arbitrary (non-zero) complex number.

In section 2. we show that the pull-back of ξ_{A-B} induced by κ (or any other function in the homotopy class of κ [7]) is isomorphic to the trivial Hopf bundle $\xi_0 : S^1 \rightarrow S^2 \times S^1 \xrightarrow{pr_1} S^2$ corresponding to zero Chern number, and therefore to zero magnetic charge in the bundle theory of magnetic monopoles. In section 3., through the composition of bundle maps, we obtain a canonical map $\xi_0 \rightarrow \xi_1$, thus establishing a tight relation between ξ_0 , the "sandwiched" bundle ξ_{A-B} , and ξ_1 , where the unit magnetic charge is described. Section 4. is devoted to a conclusion.

2 Isomorphism Between $\kappa^*(\xi_{A-B})$ and ξ_0

κ induces the pull-back bundle $\kappa^*(\xi_{A-B})$ with total space $P_{\kappa^*(\xi_{A-B})} \subset S^2 \times (\mathbb{C}^* \times S^1)$ defined by

$$P_{\kappa^*(\xi_{A-B})} = \{(b, (z, e^{i\varphi})) \mid \kappa \circ pr_1 = pr_1 \circ pr_2\} \tag{2}$$

i.e. such that the following diagram commutes:

$$\begin{array}{ccc}
 P_{\kappa^*(\xi_{A-B})} \times S^1 & \xrightarrow{pr_2 \times Id_{S^1}} & (\mathbb{C}^* \times S^1) \times S^1 \\
 \psi_{A-B}^* \downarrow & & \downarrow \psi_{A-B} \\
 P_{\kappa^*(\xi_{A-B})} & \xrightarrow{pr_2} & \mathbb{C}^* \times S^1 \\
 pr_1 \downarrow & & \downarrow pr_1 \\
 S^2 & \xrightarrow{\kappa} & \mathbb{C}^*
 \end{array}$$

Diagram 1

where ψ_{A-B} and ψ_{A-B}^* are the right actions of S^1 over $P_{\xi_{A-B}} \cong \mathbb{C}^* \times S^1$ and $P_{\kappa^*(\xi_{A-B})}$, and pr_2 is the projection in the second entry. In fact, for the lower part of the diagram one has $\kappa \circ pr_1(b, (z, e^{i\varphi}))\kappa(b) = z_0 = pr_1 \circ pr_2(b, (z, e^{i\varphi})) = pr_1(z, e^{i\varphi}) = z$, and so

$$P_{\kappa^*(\xi_{A-B})} = \{(b, (z_0, e^{i\varphi})), b \in S^2, \varphi \in [0, 2\pi)\} \tag{3}$$

while for the upper part of the diagram it holds $\psi_{A-B} \circ (pr_2 \times Id_{S^1}) = pr_2 \circ \psi_{A-B}^*$:

$$\begin{aligned}
 \psi_{A-B} \circ (pr_2 \times Id_{S^1})((b, (z, e^{i\varphi})), e^{i\lambda}) &= \psi_{A-B}(pr_2((b, (z, e^{i\varphi})), e^{i\lambda})) = \psi_{A-B}((z, e^{i\varphi}), e^{i\lambda}) = (z, e^{i(\varphi+\lambda)}), \\
 pr_2 \circ \psi_{A-B}^*((b, (z, e^{i\varphi})), e^{i\lambda}) &= pr_2(b, (z, e^{i(\varphi+\lambda)})) = (z, e^{i(\varphi+\lambda)}).
 \end{aligned}$$

Since z_0 is a fixed (but otherwise arbitrary) non-zero complex number,

$$\Phi : P_{\kappa^*(\xi_{A-B})} \rightarrow S^2 \times S^1, (b, (z_0, e^{i\varphi})) \mapsto (b, e^{i\varphi}) \tag{4}$$

is an homeomorphism, and therefore one has the bundle isomorphism

$$\xi_0 \xrightarrow{(Id_{S^2}, \Phi^{-1})} \kappa^*(\xi_{A-B}) \tag{5}$$

at the extreme left of Diagram 2.

Remark: The existence of an isomorphism $\xi_0 \rightarrow \kappa^*(\xi_{A-B})$ can be proved from the fact that $\iota \circ \kappa$ is constant and therefore the Chern class of $\kappa^*(\xi_{A-B})$ is zero.

3 Bundle Map $\xi_0 \rightarrow \xi_D$

Putting together the pull-back of ξ_D by ι , namely $\iota^*(\xi_D)$, and the isomorphism $\xi_{A-B} \xrightarrow{(Id_{\mathbb{C}^*}, \Psi^{-1})} \iota^*(\xi_D)$, respectively given in sections 4. and 5. of Ref. [1], and the results of the previous section, one obtains the sequence of bundle maps

$$\xi_0 \xrightarrow{\cong} \kappa^*(\xi_{A-B}) \longrightarrow \xi_{A-B} \xrightarrow{\cong} \iota^*(\xi_D) \longrightarrow \xi_D \tag{6}$$

represented in detail in the following commuting diagram:

$$\begin{array}{ccccccc}
 S^2 \times S^1 & \xrightarrow{\Phi^{-1}} & P_{\kappa^*(\xi_{A-B})} & \xrightarrow{pr_2} & \mathbb{C}^* \times S^1 & \xrightarrow{\Psi^{-1}} & P_{\iota^*(\xi_D)} & \xrightarrow{pr_2} & S^3 \\
 pr_1 \downarrow & & pr_1 \downarrow & & pr_1 \downarrow & & pr_1 \downarrow & & \downarrow \pi_D \\
 S^2 & \xrightarrow{Id_{S^2}} & S^2 & \xrightarrow{\kappa} & \mathbb{C}^* & \xrightarrow{Id_{\mathbb{C}^*}} & \mathbb{C}^* & \xrightarrow{\iota} & S^2
 \end{array}$$

Diagram 2

(For simplicity, we omit the upper part of the diagram, involving the right actions $\psi_0, \psi_{A-B}^*, \psi_{A-B}, \psi_D^*$ and ψ_D of S^1 on $S^2 \times S^1, P_{\kappa^*(\xi_{A-B})}, \mathbb{C}^* \times S^1, P_{\iota^*(\xi_D)}$ and S^3 , respectively.)

Since the composition of bundle maps is a bundle map, one obtains the map between the Hopf bundles ξ_0 and ξ_1 :

$$\xi_0 \xrightarrow{(\alpha, \bar{\alpha})} \xi_1 \tag{7}$$

where

$$\alpha = \iota \circ Id_{\mathbb{C}^*} \circ \kappa \circ Id_{S^2}, \alpha(b) = z_0 \tag{8}$$

and

$$\bar{\alpha} = pr_2 \circ \Psi^{-1} \circ pr_2 \circ \Phi^{-1}, \quad \bar{\alpha}(b, e^{i\varphi}) = \frac{(z_0, 1)}{\|(z_0, 1)\|} e^{i\varphi} \tag{9}$$

represented in diagram 3:

$$\begin{array}{ccc} (S^2 \times S^1) \times S^1 & \xrightarrow{\bar{\alpha} \times Id_{S^1}} & S^3 \times S^1 \\ \psi_0 \downarrow & & \downarrow \psi_D \\ S^2 \times S^1 & \xrightarrow{\bar{\alpha}} & S^3 \\ pr_1 \downarrow & & \downarrow \pi_D \\ S^2 & \xrightarrow{\alpha} & S^2 \end{array}$$

Diagram 3

4 Conclusion

Eq. (6) or *Diagram 2* clearly exhibit the fact that the bundle corresponding to the $A - B$ effect, ξ_{A-B} , is sandwiched between the Hopf bundles corresponding to zero and unit Chern numbers, respectively ξ_0 and ξ_1 . A remarkable fact of this construction is its uniqueness, in the sense that the inclusion $\iota : \mathbb{C}^* \rightarrow S^2$ is *canonical* and the map $\kappa : S^2 \rightarrow \mathbb{C}^*$ is *unique up to homotopy* (the arbitrariness of the choice of $z_0 \in \mathbb{C}^*$ is irrelevant). This establishes a strong mathematical relation between the $U(1)$ -bundles ξ_0, ξ_{A-B} and ξ_D , and therefore between the $A - B$ effect and the Dirac magnetic monopoles when these are described in the language of fiber bundle theory.

Finally, we want to comment that any smooth or continuous map $\alpha : S^2 \rightarrow S^2, \eta \rightarrow \alpha(\eta)$, induces a bundle morphism $\xi_0 \xrightarrow{(\alpha, \bar{\alpha})} \xi_1$ with $\bar{\alpha}(\eta, e^{i\varphi}) = \frac{(\alpha(\eta), 1)}{\|(\alpha(\eta), 1)\|} e^{i\varphi}$. This infinite set of morphisms accommodates into an infinite set of homotopy classes indexed by the integers since $\pi_2(S^2) \cong \mathbb{Z}$. The morphism $(\alpha, \bar{\alpha})$ given by (8) and (9) is the one induced by the $A - B$ bundle.

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